

On the relation between covariant and canonical Quantum Gravity [\[arXiv:gr-qc/1307.5885\]](https://arxiv.org/abs/1307.5885)

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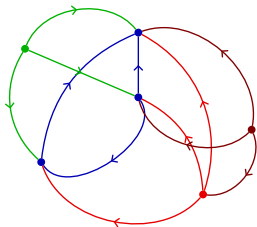
with Thomas Thiemann



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Motivation

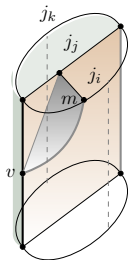
Canonical LQG



$$T_s(A) = \text{Tr}[\prod_l h_l^j(A) \prod_n t_n]$$

$$\mathcal{H}_{kin} = \bigoplus_{\gamma \in \Sigma} \mathcal{H}_{kin, \gamma}$$

Covariant LQG

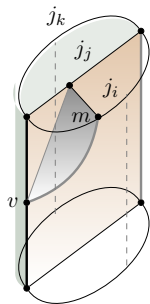


$$Z[\kappa] = \sum_c \prod_v \mathcal{A}_v \prod_f \mathcal{A}_f \times \mathcal{B}$$

$$\mathcal{H}_{\partial\kappa} \simeq \mathcal{H}_{kin, \gamma}$$

Is it possible to merge covariant and canonical LQG?

The general idea



$$\langle T_s | Z[\kappa] | T_{s'} \rangle$$
$$\partial\kappa = \gamma_s \cup \gamma_{s'}$$

Rovelli, Reisenberger, Barrett, Crane, Engle,
Pereira, Livine, Freidel, Krasnov,
Kaminski, Kisielowski, Lewandowski, . . .

Heuristic Idea

$$\delta(H) = \int \exp(itH) dt \leftrightarrow \sum_{\kappa} \underbrace{Z[\kappa]}_{\text{Feynman diagrams}}$$

[Reisenberger, Rovelli]

“Physical Scalar Product”

$$\eta[T_s](T_{s'}) := \sum_{\kappa: \gamma_{s'} \rightarrow \gamma_s} \langle T_s | Z[\kappa] | T_{s'} \rangle$$

Main Results

The extension of the [Engle,Pereira,Rovelli, Livine]-[Freidel, Krasnov]-model by [Kaminski, Kisielowski, Lewandowski] enables to investigate whether spin foams can be used to construct a rigging map for any of the presently defined Hamiltonian constraint operators.

We suggest a 'rigging map' closely following the ideas of [Reisenberger, Rovelli].

In the analysis of the resulting object we are able to identify an elementary spin foam transfer matrix that allows to generate any finite foam as a finite power of the transfer matrix.

It transpires that the postulated map does not define a projector on the physical Hilbert space. However, it might be possible to construct a proper rigging map in terms of a modified transfer matrix.

Outline

- 1 A spin foam rigging map
 - Rigging maps and spin foams - a brief review
 - Proposal for a spin foam 'rigging map'
 - Properties
- 2 Why the would-be 'rigging map' is not a proper rigging map
 - Time ordering
 - Factorization
 - Why η is not a rigging map
- 3 Discussion and outlook
 - What is going wrong?
 - Open questions

Rigging maps

How to solve the constraints of canonical QG?

Generalized solution of C

.. is a state $I \in \mathcal{D}_{kin}^*$ s.t. $[C^*I](f) := I(C^\dagger f) = 0 \quad \forall f \in \mathcal{D}_{kin}$

$\mathcal{D}_{kin} \subset \mathcal{H}_{kin}$ dense domain, algebraic dual \mathcal{D}_{kin}^*

Construction principle? What is the scalar product?

Rigging map

Given $\eta : \mathcal{H}_{kin} \rightarrow \mathcal{D}_{kin}^*$ s. t. $\langle \eta(f), \eta(f') \rangle_{phys} = [\eta(f')](f)$

Euclidean spin foam models

- 1 **Action:** $S_{BF} = \int_{\mathcal{M}} \text{Tr}[(B + \frac{1}{\beta} * B) \wedge F] + \text{Simplicity constraint}$

\mathcal{M} space-time, B bivector, F curvature of $\text{spin}(4)$ -connection on \mathcal{M} , β Barbero-Immirzi parameter

- 2 **Discretize:** \mathfrak{T} triangulation of \mathcal{M} with dual 2-complex $\kappa_{\mathfrak{T}} \rightsquigarrow S_{BF}[\kappa_{\mathfrak{T}}]$
- 3 **Quantize** $\int \mathcal{DP} e^{iS_{BF}[\kappa_{\mathfrak{T}}]} \rightsquigarrow \mathcal{A}^{BF}[\kappa_{\mathfrak{T}}, \psi_{in}, \psi_{out}] = \langle \psi_{in}, Z^{BF}[\kappa_{\mathfrak{T}}] \psi_{out} \rangle$
- 4 **Implement simplicity constraint**

Result

$$\mathcal{A}^{EPRL-FK}[\kappa_{\mathfrak{T}}, \psi_{in}, \psi_{out}] = \langle \psi_{in}, Z[\kappa_{\mathfrak{T}}] \psi_{out} \rangle =$$

$$\sum_{j_f^{\pm}, l_e^{EPRL}} \prod_f \mathcal{A}_f \prod_{e \in \kappa_{int}} Q_e \prod_{v \in \kappa_{int}} \mathcal{A}_v(j_f^{\pm}, l_e^{EPRL}) \mathcal{B}(\psi_{in}, \psi_{out})$$

$j^{\pm} = \frac{|\mathbf{1} \pm \beta|}{2} j$, l_e^{EPRL} intertwiner coupling j and (j^+, j^-)

Merging canonical and covariant LQG - Prerequisites

Can we construct $\eta[T_S](T_{S'}) = \sum_{\kappa: \gamma_{S'} \rightarrow \gamma_S} \langle T_S | Z[\kappa] | T_{S'} \rangle$?

Need identification of boundary states ψ_{in}, ψ_{out} with states in \mathcal{H}_{kin}

Canonical LQG

Continuum theory

Quantize point separating
sub-algebra of Poisson algebra

\rightsquigarrow all possible graphs in Σ

Hamiltonian \rightsquigarrow 3-valent nodes

EPRL-FK Model (Eucl.)

Discrete theory

κ dual to triangulation \mathfrak{T} of \mathcal{M} ,

\rightsquigarrow boundary graphs dual to
triangulation of Σ

at least 4-valent nodes

Other issues:

Path integral measure? Fate of observables?

Implementation of simplicity constraint?

Abstract foams - The KKL-model

Spin foams - a tool to determine the physical scalar product ?

KKL-model

[Kaminski, Kisielowski, Lewandowski]

Achievements:

- 1 $Z[\kappa]$ defined for arbitrary 2-complexes allowing arbitrary boundary graphs $\partial\kappa$
- 2 For certain choice of β : $\mathcal{H}_{\partial\kappa} \simeq \mathcal{H}_{kin,\partial\kappa}$

In the following we will consider **abstract foams**

Abstract foam κ : p.l. homeomorphic to a 2-complex $\{f, e, v\}$ whose boundary $\partial\kappa$ is the disjoint union of **closed graphs bordering κ**

γ borders κ : exist 1-to-1 affine map $\gamma \times [0, 1] \rightarrow \kappa$

Proposed rigging-map

Abstract equivalence

Two embedded spin nets belong to the same abstract equivalence class $[s]_A$ if they are embeddings of the same abstract spin net s_A .

A spin foam rigging map

$$\eta[T_s](T_{s'}) = \sum_{[s']_A \in N_A} \eta_{[s]_A, [s']_A} L_{[s']_A} \quad \text{with} \quad \eta_{[s]_A, [s']_A} = \sum_{\kappa_A(s_A, s'_A)} Z[\kappa_A(s_A, s'_A)]$$

$$\text{and} \quad L_{[s']_A} = \eta_{[s']_A} \sum_{\hat{s} \in [s']_A} \langle T_{\hat{s}}, \cdot \rangle$$

$$Z[\kappa_A(s_A, s'_A)] := \sum_{j_f, \ell_e} \prod_f \mathcal{A}_f \prod_{e \in \kappa_{int}} Q_e \prod_{v \in \kappa_{int}} \mathcal{A}_v(j_f, \ell_e) \prod_{l \in \partial \kappa} \delta_{j_l, j_{f_l}} \prod_{n \in \partial \kappa} \delta_{\ell_n, \ell_{e_n}}$$

[Kaminski, Kisielowski, Lewandowski], [Ding, Han, Rovelli], [Bahr, Hellmann, Kaminski, Kisielowski, Lewandowski]

Properties of $Z[\kappa]$

$$Z[\kappa_A(s_A, s'_A)] := \sum_{j_f, \ell_e} \prod_f \mathcal{A}_f \prod_{e \in \kappa_{int}} Q_e \prod_{v \in \kappa_{int}} \mathcal{A}_v(j_f, \ell_e) \prod_{l \in \partial \kappa} \delta_{j_l, j_{f_l}} \prod_{n \in \partial \kappa} \delta_{\ell_n, \ell_{e_n}}$$

[Kaminski, Kisielowski, Lewandowski], [Ding, Han, Rovelli], [Bahr, Hellmann, Kaminski, Kisielowski, Lewandowski]

Can be defined s.t. $Z[\kappa_A(s_A, s'_A)]$ is invariant if ...

- ... faces labeled by trivial representation are added or removed
- ... internal edges splitting a face are added or removed
- ... 2-valent vertices are added or removed

$Z[\kappa_A](s_A, s'_A) : \mathcal{H}_{kin, \gamma}(s'_A) \rightarrow \mathcal{H}_{kin, \gamma}(s_A)$ is cylindrically consistent!

Furthermore: $\forall \gamma \exists \kappa_\gamma^0$ s.t. $Z[\kappa_\gamma^0] : \mathcal{H}_{kin, \gamma} \rightarrow \mathcal{H}_{kin, \gamma}$ with $Z[\kappa_\gamma^0] = \mathbb{1}_\gamma$

Minimal foams

Note

Since adding faces with $j_f = 0$, splitting faces and edges leave the amplitude/spin net function invariant one should only sum over equivalence classes in the rigging map .

Definition

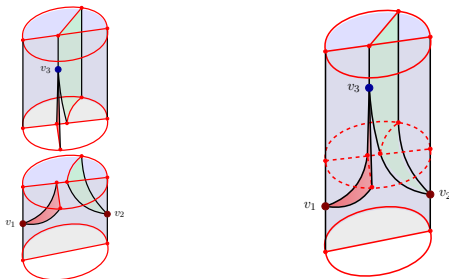
An abstract foam/graph is called **minimal** iff it cannot be obtained from another foam/graph by subdivisions.

Remark

- Minimal representative is not unique . . .
- . . . but amplitude does not depend on choice

Can always fix a representative for each class and only sum over those.

Gluing



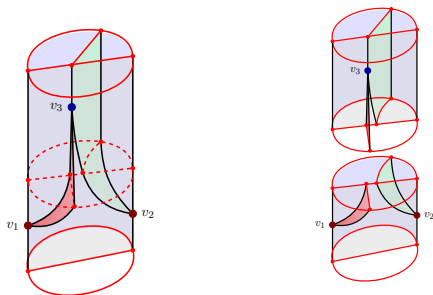
Suppose $\kappa_1 \cap \kappa_2 = \partial\kappa_1 \cap \partial\kappa_2 = \tilde{\gamma}$ then

$$\sum_{\tilde{s}(\tilde{\gamma})} Z[\kappa_1(s, \tilde{s})] Z[\kappa_2(\tilde{s}, s')] = Z[\kappa_1 \# \kappa_2(s, s')]$$

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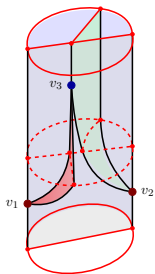
Splitting



Idea: Use this to split big complexes to get a better control on

$$\eta[T_s](T_{s'}) := \sum_{\kappa: \gamma_{s'} \rightarrow \gamma_s} \langle T_s | Z[\kappa] | T_{s'} \rangle$$

Time ordering



Definition

- $v \in \kappa_{int}$ s.t. $\exists n \in \gamma_i$ and $\exists e \in \kappa_{int}$ with $s(e) = n$ and $t(e) = v$
 \leadsto **vertex of first generation**
- Inductively: Vertex of n^{th} generation
- Only final graph \Rightarrow count backwards
- $\partial\kappa = \emptyset$ all $v \in \kappa$ of first generation

Theorem

A minimal foam κ can be **uniquely split** into minimal foams κ_i , called '**one time step**' foams, containing **only vertices of i^{th} generation** with respect to the original foam $\kappa = \kappa_1 \# \cdots \# \kappa_N$ where N is the **maximal generation** of κ .

Sketch of the Proof I

To prove the theorem we need the following observations:

Lemma

If $e \in \kappa_{int}$ has a vertex of n^{th} generation then the other vertex of e is either of generation $n - 1, n$ or $n + 1$ or a vertex in a final graph.

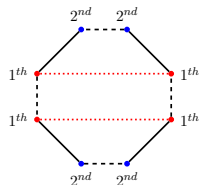
Introduce **additional vertices** so that the **second vertex of e** is either of **generation $n - 1, n$ or $n + 1$** .

Lemma

If a face has an edge e joining a vertex of n^{th} and $(n + 1)^{th}$ generation then it contains at least one other edge connecting vertices of n^{th} and $(n + 1)^{th}$ generation (or v of generation m and v' in final graph).

Sketch of the Proof II

Proof of the theorem



- $\mathcal{E}_{n,n+1} := \{e \in \kappa_{int} : \text{joining vertices of } n^{th} \text{ and } (n+1)^{th} \text{ generation}\}$
- Number of boundary edges $e \in \mathcal{E}_{n,n+1}$ of f is even
- Split faces until a face contains at most 2 edges in $\mathcal{E}_{1,2}$
- Split all edges in $\mathcal{E}_{1,2}$ and join new vertices by edges

$$\leadsto \kappa = \kappa_1 \# \kappa'$$

Repeat the procedure until $\kappa = \kappa_1 \# \dots \# \kappa_N$

Factorization of the rigging map I

If η is a proper rigging map then ...

$$\eta[T_{s_f}](\hat{H}T_{s_i}) \stackrel{!}{=} 0 \quad \forall T_{s_i} \in \mathcal{H}_{kin}$$

$$\implies \eta[T_{\emptyset}](\hat{H}T_{s_i}) = \left(\sum_n (\eta^cT_{\emptyset})^n \right) \eta^c[T_{\emptyset}](\hat{H}T_{s_i}) \stackrel{!}{=} 0$$

where \sum_{κ} in η^c is restricted to connected foams

Thus to test the rigging map it suffices to only consider connected foams!!

- But when splitting a connected κ into 'one time step' foams $\kappa_1 \# \dots \# \kappa_n = \kappa$ then κ_i are not necessarily connected
- However, if T_{\emptyset} is excluded and either $\gamma(s_f)$ and or $\gamma(s_i)$ are connected then any foam $\kappa_1 \# \dots \# \kappa_n$ is connected

Test whether $\eta^c[T_{s_f}]$ annihilates H for any s_f with connected graph $\gamma(s_f)$

Factorization of the rigging map II

Spin foam transfer matrix

$\hat{K}_{\gamma, \gamma'}$ set of 'one time step' foam and $P_\gamma : \mathcal{H}_{kin} \rightarrow \mathcal{H}_{kin, \gamma}$

$$\hat{Z} := \sum_{\gamma, \gamma'} P_{\gamma'} \left[\sum_{\hat{\kappa} \in \hat{K}_{\gamma, \gamma'}} Z(\hat{\kappa}) \right] P_\gamma$$

We then obtain for $\gamma(s_f)$ connected ...

$$\eta^c[T_{s_f}](T_{s_i}) = \sum_{\kappa \in K_{\gamma(s_i), \gamma(s_f)}} \langle T_{s_f}, Z(\kappa_1 \# \dots \# \kappa_N) T_{s_i} \rangle = \sum_{N=0}^{\infty} \langle T_{s_f}, \hat{Z}^N T_{s_i} \rangle$$

Remarks

- Above factorization would lead to ordering ambiguities for disconnected foams
- \hat{Z} is symmetric

A naive argument

If η would be a proper rigging map then

$$\eta^c[T_{S_f}](HT_{S_i}) = \langle T_{S_f}, \sum_{N=0}^{\infty} \hat{Z}^N HT_{S_i} \rangle \stackrel{!}{=} 0$$

But ...

$$A := \sum_{N=0}^{\infty} Z^N \Rightarrow A = \mathbb{1} + ZA$$

leads to a contradiction:

$$0 = AH = (\mathbb{1} + ZA)H = H$$

Of course, this expression is only formal as A is very likely diverging and therefore requires a regularization

Regularization of \hat{Z}

Cut-Off: $K_{\gamma, \gamma'}$ is infinite, therefore introduce cut-off (J, N_f, N_e)

Weight: Suppose there exist a weight $\omega(\kappa \parallel \kappa') = \omega(\kappa)\omega(\kappa')$ s.t.

$$Z' := \sum_{\hat{\kappa} \in \hat{K}_{\gamma, \gamma'}} \omega(\hat{\kappa}) Z[\hat{\kappa}] \text{ has finite norm}$$

Even if Z' is now densely defined $(Z')^n$ is not necessarily densely defined.

Recall: Z is symmetric and so is Z'

Suppose: Z' can be extended to an self-adjoint operator with projection valued measure E and let

Then: Z' acts on the states $\psi_q := \int_{-q}^q dE(\lambda) \psi$, $\psi \in \mathcal{H}_{kin}$, by multiplication with $\lambda \in [-q, q]$ where $0 < q < 1$

$$\text{But: } A\psi_q = \int_{-q}^q dE(\lambda) (1 - \lambda)^{-1} \psi \rightsquigarrow A = \mathbb{1} + ZA \not\checkmark$$

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What is going wrong?

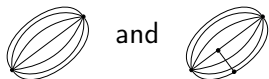
- Option I: Z is itself a projector, e.g. BF-theory
- Option II: Vertex amplitude too local
- Option III: Necessity to restrict to one Plebanski sector?
- Option IV: Wrong assumption on the weight!

What is going wrong?

Option I: Z is itself a projector, e.g. BF-theory

- ① It can be shown that for $\beta = 1$ the operator $\sum_{\kappa \in K_{\text{Melon}}} Z[\kappa]$ annihilates the Euclidean constraint. Here, $\kappa \in K_{\text{Melon}}$ are foams with just one internal

vertex and boundary graphs of the form:



[Alesci, Thiemann, A.Z.]

- ② Z can be interpreted as an object obtained from some sort of coarse graining. In [Dittrich, Hellmann, Kaminski] a similar object was derived for holonomy spin foam models via coarse graining. For BF-theory this already incorporates the dynamics.

What is going wrong?

Option II: Vertex amplitude too local

The would-be rigging map can be factorized since foams can be split, that is, $\kappa = \kappa_1 \# \cdots \# \kappa_N$

Suppose $\mathcal{A}_V Q_e(v, v') \mathcal{A}_{V'} \not\approx \mathcal{A}_V Q_e(v, m) \mathcal{A}_m Q_e(m, v') \mathcal{A}_{V'}$
then splitting is no longer possible.

In this sense: **vertex amplitude too local**

Other hints into this direction: E.g. [Dittrich], [Hellmann, Kaminski], ...

What is going wrong?

Option III: Necessity to restrict to one Plebanski sector?

$$\begin{aligned}
 2\pi\delta(C) &= \lim_{T \rightarrow \infty} \int_{-T}^T e^{itC} = \lim_{T \rightarrow \infty} \int_0^T [e^{itC} + e^{-itC}] \\
 &= \lim_{T \rightarrow \infty} \lim_{n \rightarrow \infty} \sum_{k=1}^n \frac{T}{n} \underbrace{[e^{iCT/n}]^k}_{U^k} + \underbrace{[e^{-iCT/n}]^k}_{(U^\dagger)^k}
 \end{aligned}$$

But: \hat{Z} is symmetric $\leadsto \hat{Z} = U + U^\dagger$

$$\Rightarrow \sum_{k \in \mathbb{N}} Z^k = \sum_{k \in \mathbb{N}} (U + U^\dagger)^k \neq \delta(H)$$

A similar problem occurs in the **asymptotic expansion**.

Can it be solved by the proper vertex proposal [Engle]?

What is going wrong?

Option IV: Wrong assumption on the weight!

We assumed that Z can be regularized by introducing a weight ω that obeys $\omega(\kappa_1)\omega(\kappa_2) = \omega(\kappa_1 \# \kappa_2)$. Is that sensible?

- 1 Regularization might require symmetry factors that destroy splitting:

$$\delta(x) = \int dk e^{ikx} = \int dk \left(\sum_n \frac{i^n}{n!} (kx)^n \right)$$

For example in **standard GFT** $\frac{\lambda^{|\gamma^{(0)}|}}{|\text{sym}(\gamma)|}$

- 2 Spin foam amplitudes suffer from **bubble divergencies**. However, if foams with bubbles are excluded then gluing might no longer be possible, since even if κ_1, κ_2 contain **no bubbles**, $\kappa_1 \# \kappa_2$ might have **bubbles**.

Open questions

Regularization

- Explore regularization methods of generalized GFT [Oriti]
- Explore latest results on regularization of the EPRL-model [Bonzom, Carrozza, Oriti, Puchta, Riello, Rovelli, Smerlak, ...]
- Connection to coarse graining [Dittrich, ...]?

Proper vertex

Generalize proper vertex proposal [Engle] to n-valent vertices and test whether the modified rigging map still has the same problems.

Technical aspects

- Single time step \rightsquigarrow better control on the sum
- Formalism to design single time step foams; [Kisielowski, Lewandowski, Puchta]

Thank you for your attention!

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