

Black holes and reversibility

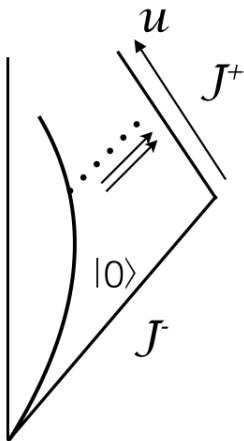
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Setup: gravitational collapse



- ▶ A **black hole** forms from ingoing matter.
- ▶ Trapping horizon forms and **peels off** outgoing geodesics.
- ▶ Thermal **Hawking radiation** is emitted.
- ▶ Breakdown of predictability?

Information loss as a physical problem

“Information loss violates a basic tenet of quantum mechanics.”

- ▶ **Information loss** happens all the time:
 - ▶ with open systems (decoherence)
 - ▶ with non-Cauchy “out” surfaces

[Wald 13]

- ▶ Information loss does **not** mess up with conservation laws.

[Banks, Peskin, Susskind 84; Unruh, Wald 95]

- ▶ The only real question is: what **difference** would it make?

What **physical effects** relate to the information loss problem?

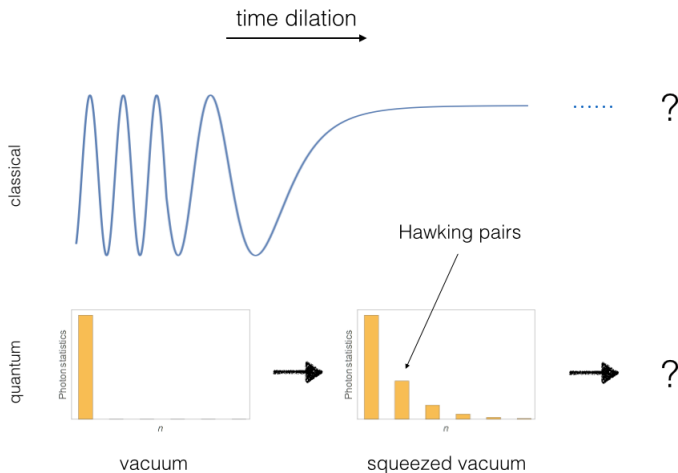
Outline

Black holes as squeezers

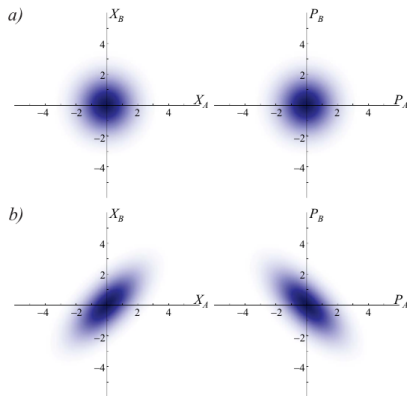
Past/future entanglement

(A)cyclic processes

The Hawking effect



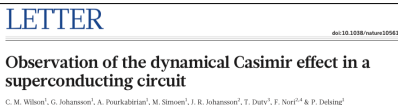
Two-mode squeezed vacuum



$$|\psi_{AB}\rangle \propto \sum_{n=0}^{\infty} (\tanh r)^n |n, n\rangle \quad \Longrightarrow \quad \rho_A \propto \sum_{n=0}^{\infty} (\tanh r)^n |n\rangle \langle n|$$

Stronger **squeezing**, higher **temperature** ($e^{-\hbar\omega/kT} = \tanh r$).

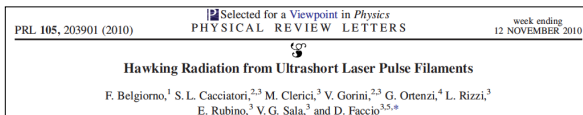
Observing TMSV



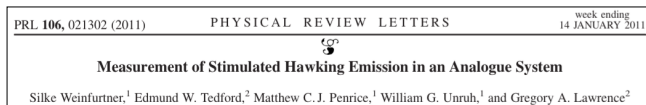
SQUID



BEC



Nonlinear optics



Hydrodynamics

Open questions

The evaporation problem is a **runaway** problem

radiation \implies mass loss \implies smaller hole \implies
higher squeezing \implies more radiation...

The questions for us are

- ▶ does this lead to an **explosive** behavior?
- ▶ does **thermality** break down at late times?
- ▶ what **astrophysical signatures** should we look for?

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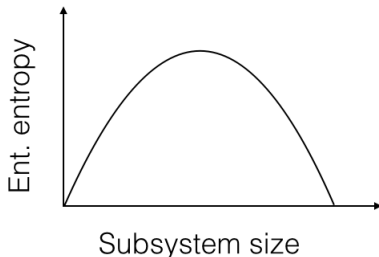
Entanglement in finite systems

In **finite dimensions**, entanglement entropy

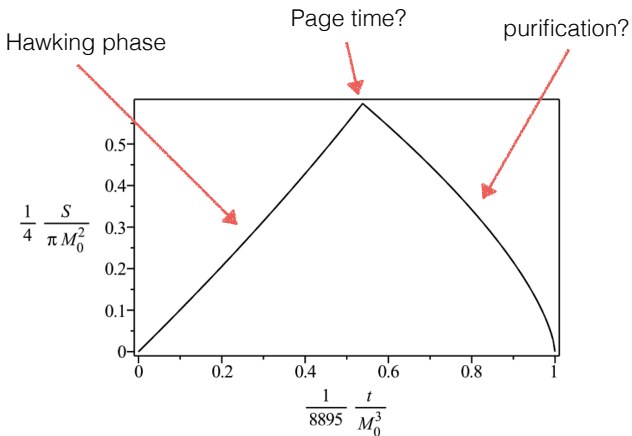
$$S[\rho_A] \equiv -\text{tr}_A[\rho_A \ln \rho_A] \quad \text{with} \quad \rho_A \equiv \text{tr}_B[\rho_{AB}]$$

is **unitarily invariant** and satisfies the **triangle inequality**

$$|S[\rho_A] - S[\rho_B]| \leq S[\rho_{AB}] \leq S[\rho_A] + S[\rho_B].$$



Page's conjecture



[Page (93,13)]

Reversibility : three open questions

Is the evaporation process

1. **unitary**, viz. is purity preserved ?

$$S_{\text{vN}}[\rho_{\text{out}}] = S_{\text{vN}}[\rho_{\text{in}}] ?$$

[Hawking (76)]

2. **cyclic**, viz. does entanglement return to its initial value?

$$\lim_{u \rightarrow +\infty} S_{\text{P}}(u) = \lim_{u \rightarrow -\infty} S_{\text{P}}(u) ?$$

[Page (93)]

3. **conservative**, viz. do energy input and output match ?

$$\lim_{u \rightarrow +\infty} M(u) = 0 ?$$

Working assumptions

Neglect

- ▶ angular momentum (of spacetime and fields)
- ▶ backscattering
- ▶ non-conformal interactions

but not

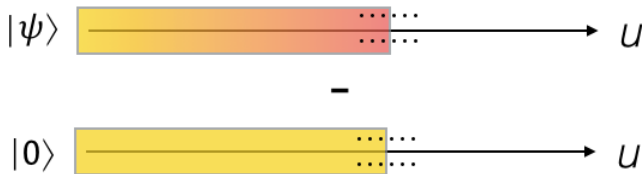
- ▶ semiclassical backreaction (even strong).

Reduces field dynamics to **2d CFT**:

$$\phi(t, r) = r^2 \int_{S^2} d\Omega^2 \Phi(t, r, \Omega)$$

Renormalized entanglement entropy

In QFT, entanglement entropy is UV-divergent. Subtract vacuum contribution



Defines **renormalized** entanglement entropy

$$S_P(u) = [\rho_\psi(u)] - S[\rho_0(u)]$$

[Holzhey, Larsen, Wilczek (94)]

The Page curve

Starting from the (non-covariant) CFT formula for a segment

$$S[\rho(R)] = \frac{1}{3} \log \frac{L(R)}{\epsilon}$$

[Holzhey, Larsen, Wilczek (94)]

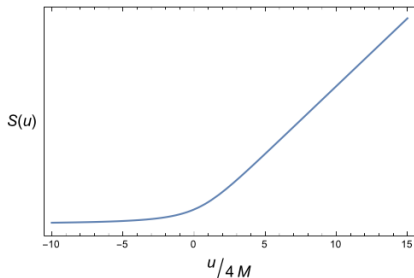
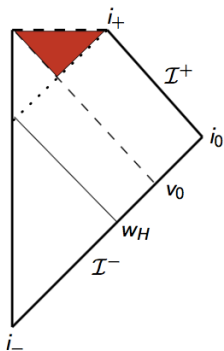
we obtain the **geometric** formula

$$S(u) = \frac{1}{12} \ln \chi(u)$$

[Bianchi, MS 14]

with $\chi = \omega_+/\omega_-$ the in-out **redshift** factor.

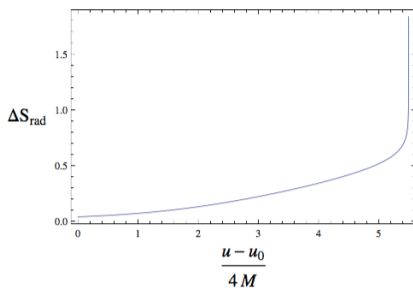
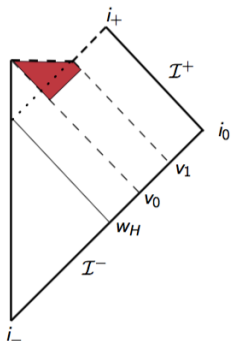
Vaidya spacetime: the Hawking phase



$$S(u) = \frac{1}{12} \log \left(\frac{1 + W(e^{-u/4M})}{W(e^{-u/4M})} \right) \sim \frac{u}{48M}$$

[Bianchi, de Lorenzo, MS 14]

“Hawking spacetime”: thunderbolt



$$S(u) \sim \frac{1}{12} \log \left(\frac{4M}{u - u_H} \right)$$

[Bianchi, de Lorenzo, MS 14]

From spacetime to the Page curve

More examples illustrate the connection between geometry and entanglement...

[Bianchi, de Lorenzo, MS 14]

... but in this approach, where

spacetime \implies entropy,

backreaction is an **input**. Next best thing after blind guess!

Importance of **other**, less narrow **approach**.

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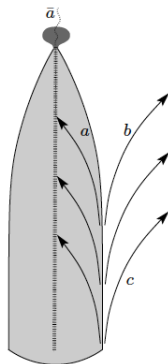
Past/future entanglement

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Unitarity violations?

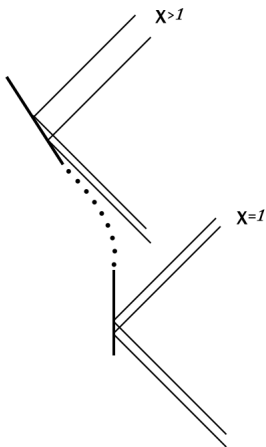
Several authors propose that evaporation is **non-unitary** (in the QFT sector):

- ▶ decoherence without dissipation: **spin bath** model
[Unruh, Wald 95; Unruh 12]
- ▶ quantum gravity decoherence: **defects** in spacetime weave
[Perez 14]



Here I'll explore another possibility: **unitary but acyclic** evaporation.

The moving mirror



- ▶ Mirror starts at rest...
- ▶ ... then accelerates...
- ▶ ... then is inertial again.

$$\Delta S \propto (\text{relative rapidity})$$

Unitary but acyclic.

What does cyclicity imply?

Outgoing energy flux

Other natural observable at \mathcal{I}^+ : **energy flux**

$$F(u) \equiv 4\pi r^2 \langle \text{in} | T_{uu} | \text{in} \rangle$$

and **Bondi mass**

$$M(u) \equiv M_0 - \int_{-\infty}^u du' F(u').$$

In the 2d approximation,

$$F(u) = -\frac{1}{24\pi} \left(\frac{\ddot{\dot{p}}(u)}{\dot{p}(u)} - \frac{3\ddot{p}(u)^2}{2\dot{p}(u)^2} \right)$$

[Fulling, Davies, Unruh (76)]

The it from bit equation

$$2\pi F(u) = 6\dot{S}(u)^2 + \ddot{S}(u)$$

- ▶ “Page curve” $S(u)$ determines energy flux $F(u)$
- ▶ Energy flux $F(u)$ determines Page curve $S(u)$, via

$$-\ddot{\psi}(u) + 12\pi F(u)\psi(u) = 0 \quad \text{where} \quad \psi \equiv e^{6S}$$

- ▶ Flux $F(u)$ is “exceptional”: $F(u) + \delta F(u)$ not a flux
- ▶ Implies quantum inequality: $|F|\tau^2 \lesssim 1$

It-from-bit and the GSL

$$2\pi F(u) = 6\dot{S}(u)^2 + \ddot{S}(u)$$

Generalizes GSL in two ways:

- ▶ Includes **non-adiabatic term** (identity rather than ineq.)
- ▶ Does not require special causal structure (event horizon)
- ▶ **Gives back GSL** when $|\ddot{S}| \ll \dot{S}^2$. For a Schwarzschild black hole, with

$$\dot{S} = \frac{1}{48M_B} \quad \text{and} \quad F = -\dot{M}_B = -\frac{\dot{S}_{\text{BH}}}{32\pi M_B}$$

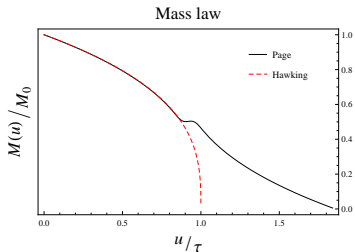
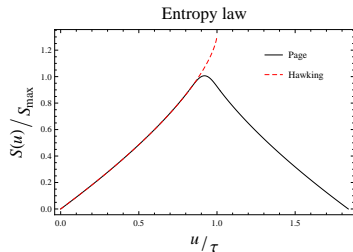
you get

$$dS_{\text{BH}} + dS = \frac{u}{96M} > 0.$$

A black hole's last gasp

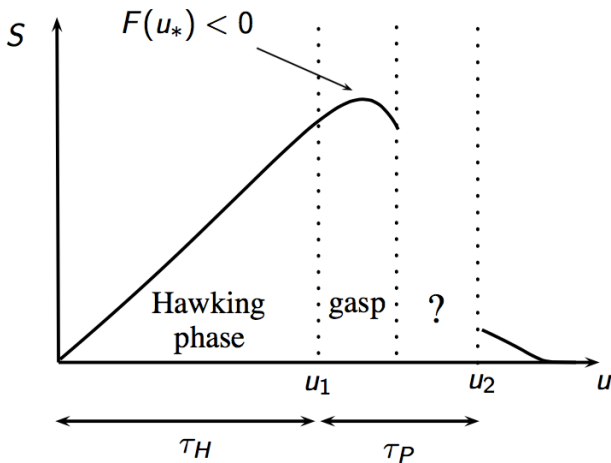
$$2\pi F(u) = 6\dot{S}(u)^2 + \ddot{S}(u)$$

At the “Page time” u_* , the flux is **negative**: $F(u_*) < 0$.



Black hole's “**last gasp**”.

Time scales



Lifetime of a black hole

From the it-from-bit equation we get that if

- ▶ the evaporation process is cyclic
- ▶ energy is conserved: $M_B(u) > 0$,

then the **purification time** must be **large**:

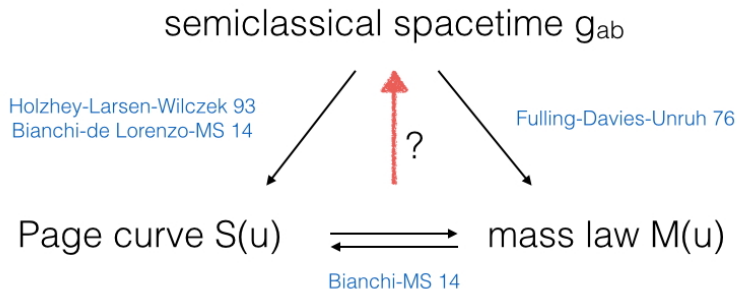
$$\tau_P \geq \xi \frac{(M_0^2 - M_1^2)^2}{M_1 m_P^2} = \begin{cases} \mathcal{O}(M_0^4/m_P^3) & \text{if } M_1 = \mathcal{O}(m_P) \\ \mathcal{O}(M_0^3/m_P^2) & \text{if } M_1 = \mathcal{O}(M_0/2) \end{cases}$$

[Carlitz-Willey 87; Bianchi, MS 19]

Recent nonsingular black hole models fail to respect this bound.

[Frolov, Vilkoviski 91; Hayward 06; Bardeen 14; Rovelli, Vidotto, Haggard 14]

An open problem



1. Is there a nonsingular black hole spacetime such that evaporation is **cyclic** and (sub)-**conservative**?
2. What kind of spacetime does Page's curve describe?

Conclusions

- ▶ Focus on **asymptotic observers** (us).
- ▶ In field theory, **unitarity** is **not** equivalent to **cyclicity**.
- ▶ From a (guessed) geometry, can **compute** the Page curve.
- ▶ **Inverse problem** seems insightful, thanks to **it-from-bit**.

Thanks to

A. Ashtekar, E. Bianchi, T. de Lorenzo, A. Perez, C. Rovelli.